

Tracing fingerprints of cosmic evolution in the large scale structures of the Universe

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October 15, 2019

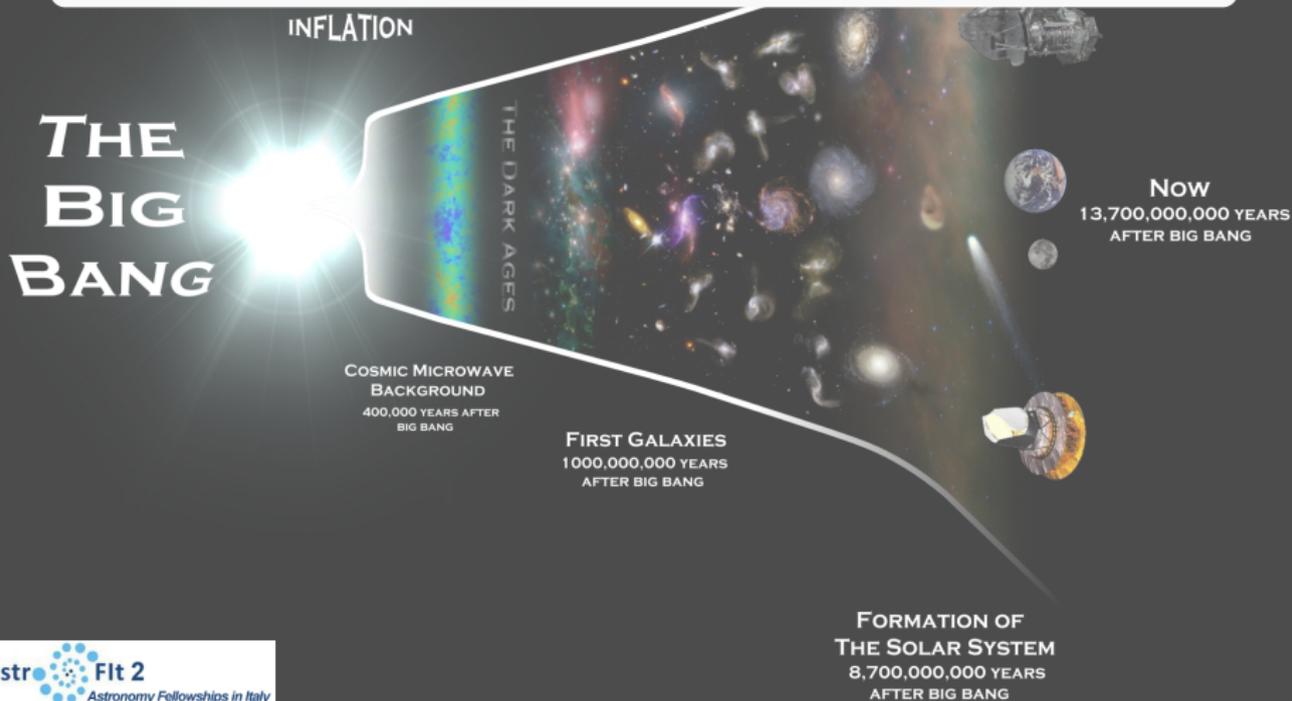
Main epochs of discussion



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Inflation

Initial rapid expansion of the Universe, generation of primordial fluctuations

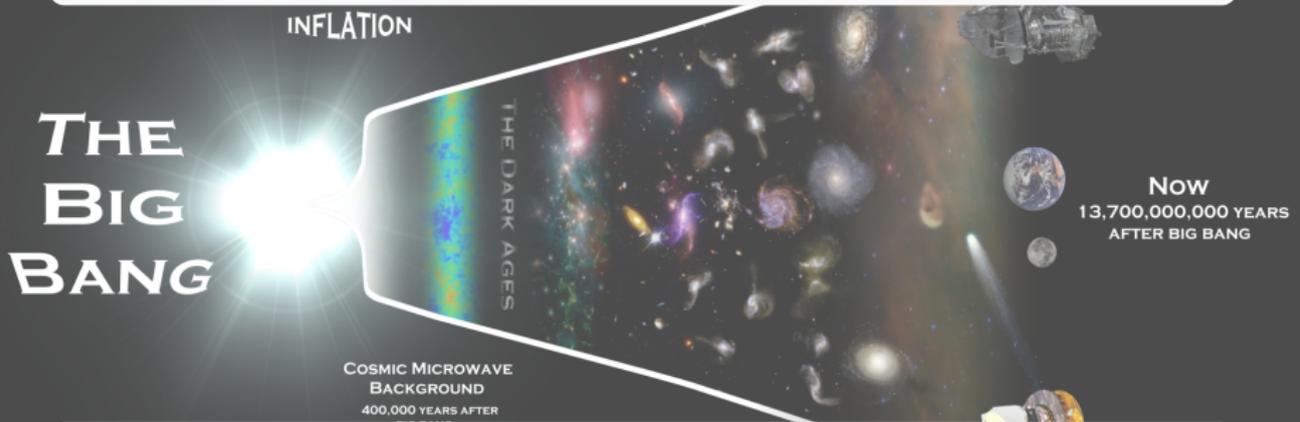


Main epochs of discussion



Inflation

Initial rapid expansion of the Universe, generation of primordial fluctuations



Reionization

During matter dominated phase, the process of ionizing hydrogen, helium atoms in the intergalactic medium by high energy photons

FORMATION OF
THE SOLAR SYSTEM
8,700,000,000 YEARS
AFTER BIG BANG



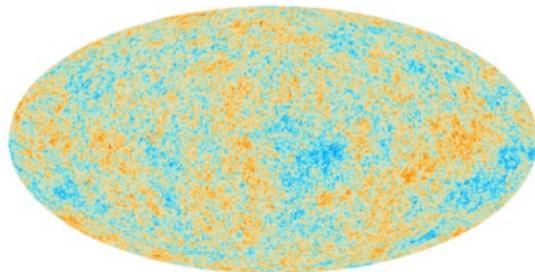
Reconstructing primordial Universe

CMB anisotropy map \implies Angular power spectrum (C_ℓ^T)

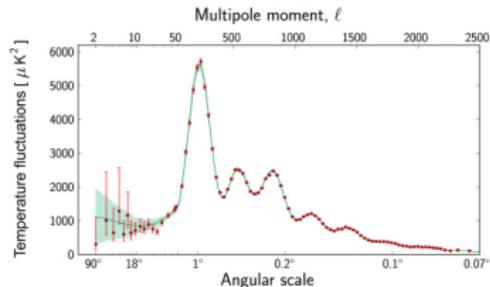


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Planck CMB map



Angular power spectrum (Planck)



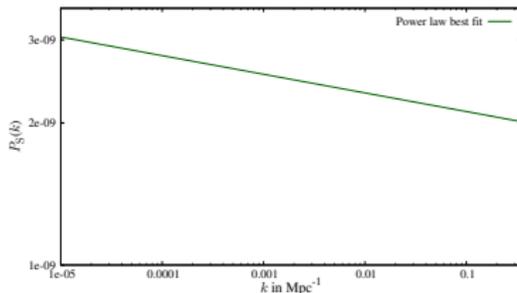
From the map of temperature anisotropy, we obtain the its power spectrum in angular scales

Primordial power spectrum (P_k) \implies Angular power spectrum (C_ℓ^T)

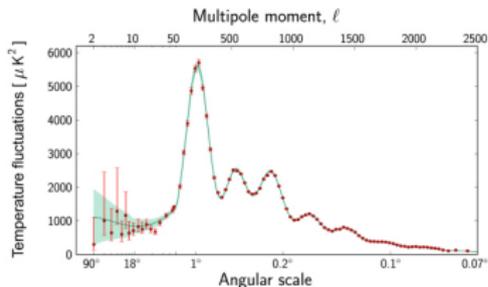


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Primordial power spectrum



Angular power spectrum (Planck)



$$C_\ell^T = \sum_i G_{\ell k_i} P_{k_i}$$

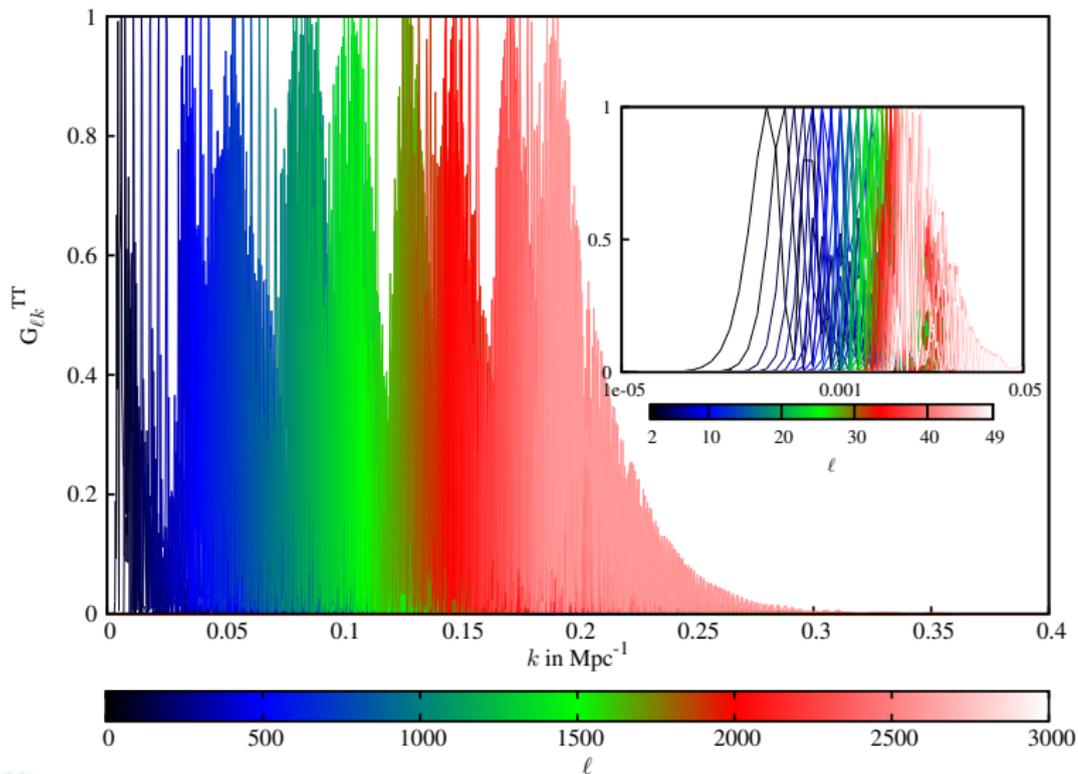
$G_{\ell k}$ is the radiative transport kernel.

Transport kernel ($G_{\ell k}$)



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Transport kernel for temperature anisotropy computed using CAMB



To search for the shape of primordial spectra:

I use 2 approaches:

- The reverse engineering or top down approach:

Directly from the data (say CMB anisotropy), I use different **deconvolution algorithms** to trace back or reconstruct the primordial signal

- Model building and fitting or bottom up approach:

Using proper theoretical/phenomenological models of inflation, I confront them with the data

Both the approaches have to be used efficiently in order **to extract the most out of the data**. Such as, the reconstruction provides the hints for building models in the shape of the data

Richardson-Lucy algorithm : Origin



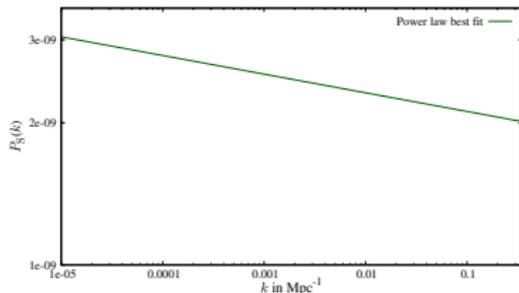
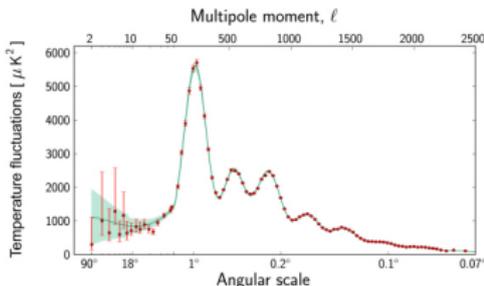
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This deconvolution iteratively solves for the PPS:

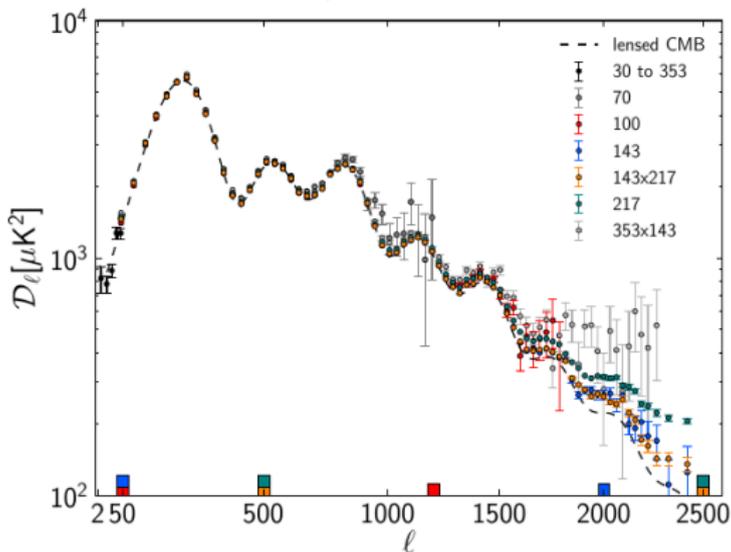
Richardson (1972) and Lucy (1974)

$$P_k^{(i+1)} - P_k^{(i)} = P_k^{(i)} \times \left[\sum_{\ell} \tilde{G}_{\ell k} \left(\frac{C_{\ell}^D - C_{\ell}^{T(i)}}{C_{\ell}^{T(i)}} \right) \right]$$

PPS at $i + 1$ 'th iteration is obtained as a correction factor to the i 'th PPS through the deconvolution



Angular power spectra (in different Planck frequencies)



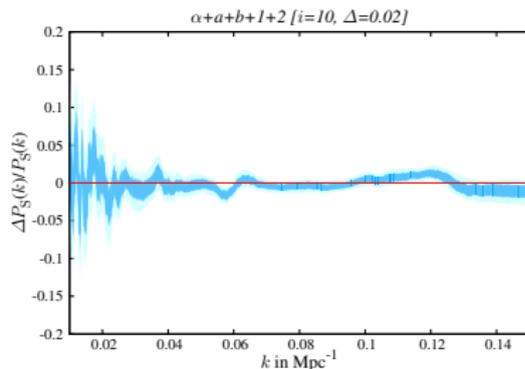
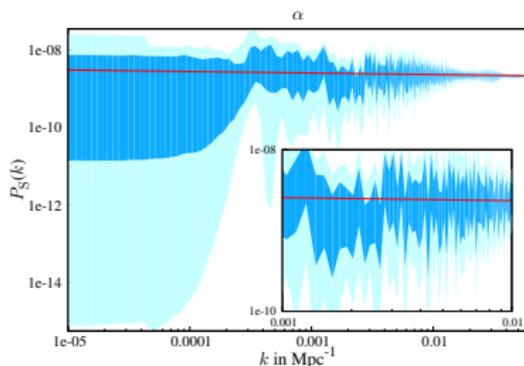
- 5 different spectra for parameter estimation, calculated from combinations of maps in different frequency channels
- Foreground and calibration effects
- Substantial lensing

Features and error estimation

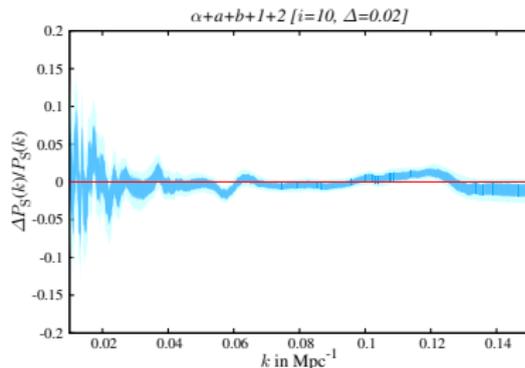
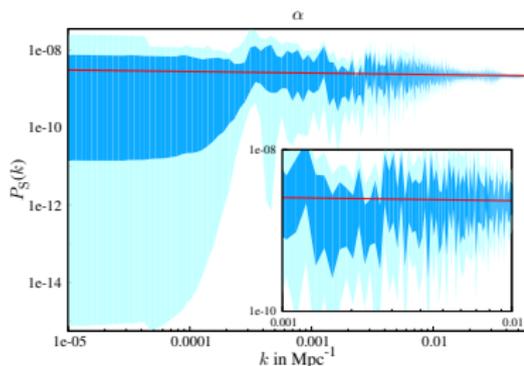


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We generate and use 1000 mock angular power spectra in every Planck channels. Using 1000 reconstructed spectra we obtain the 68% and 95% confidence contours



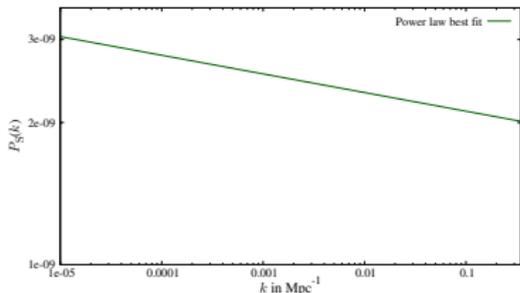
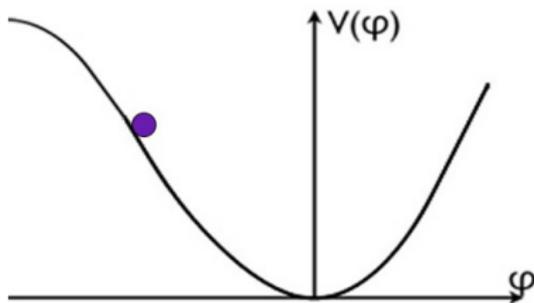
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Standard model assumption, power law is consistent at all scales apart from few localized oscillations, near $\ell \simeq 22, 200 - 300, 750 - 850$

Hazra, Shafieloo and Souradeep, JCAP 2014

Inflation is rapid expansion of the Universe, driven by presumably a scalar field, rolling slowly down its nearly flat potential



Scalar perturbations generated from the fluctuations in the field does not have strong scale dependence



In order to generate the features indicated by the reconstruction, we need departure from the strict slow roll inflaton

Start from a **faster roll potential** and **transition to a slow roll potential** or a step in the inflaton potential generates the required features.

Extensive works on feature model have been carried out **We are not the only ones** : Starobinsky-1992; Linde, Sasaki, Tanaka-1999; Adams, Creswell, Easter-2001; Covi, Hamann, Melchiorri, Slosar, Sorbera-2006; Joy, Sahni, Starobinsky-2008; Jain, Chingangbam, Gong, Sriramkumar, Souradeep-2009; Mortonson, Dvorkin, Peiris, Hu-2009; Hazra, Aich, Jain, Sriramkumar, Souradeep-2010; Flauger, McAllister, Pajer, Westphal and Xu-2010; Aich, Hazra, Sriramkumar, Souradeep-2011; Bousso, Harlow, Senatore-2013; Meerburg, Spergel, Wandelt-2014
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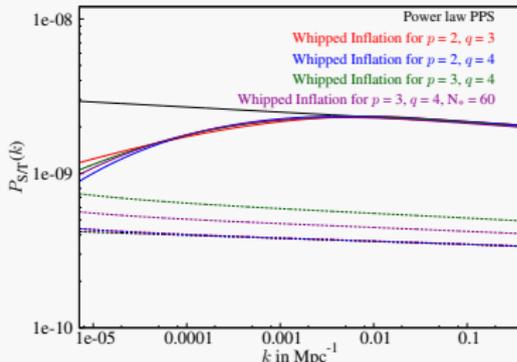
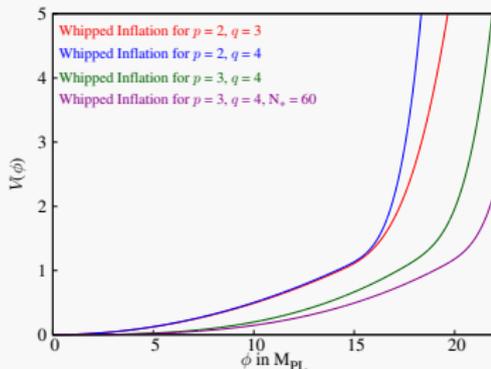
We were the first to generate a wide class of features in a single framework of inflaton potential

Hazra, Shafieloo, Smoot, Starobinsky, PRL (2014) [Editor's suggestion]

Whipped Inflation potential

$$V(\phi) = V_S(\phi) + \gamma V_R(\phi)$$

Moderate fast-roll \implies strict slow-roll



In a continuous potential Whipped Inflation provides **large scale scalar suppression** (without a running like small scale suppression), **low non-Gaussianity**. The tensor amplitude depends on the scale of $V(S)$.

BINGO: One code to solve them



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Based on **Maldacena 2002** formalism, see **Chen et al. 2006-2010** for initial works on feature model non-Gaussianities

BINGO : BI-spectra and Non-Gaussianity Operator, **Hazra, Sriramkumar and Martin, JCAP 2013**; **Sreenath, Hazra, Sriramkumar, JCAP 2015**

- BINGO solves the scalar field equation for background and the curvature perturbation equation (Mukhanov-Sasaki equation) for any canonical scalar field driven inflaton
- Calculates the bispectra for arbitrary triangular configurations [**First public code to calculate the bispectrum from inflation**]
- Performs complete numerical calculation without any slow roll approximations (MPI enabled)
- Easy to fuse with CAMB and COSMOMC for parameter estimation

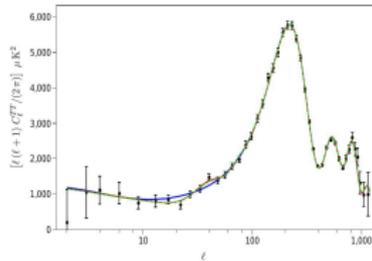
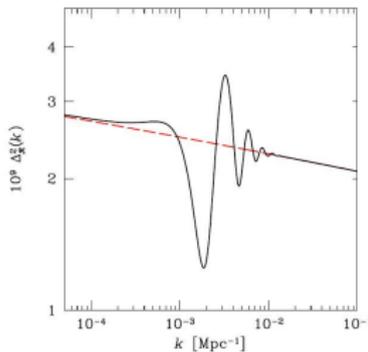
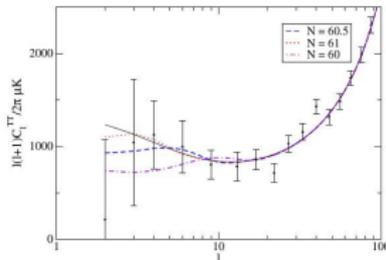
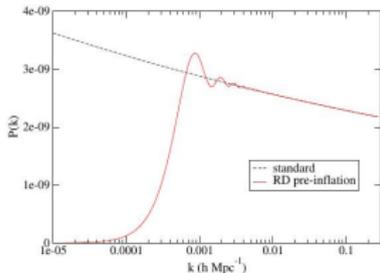
Features : earlier works



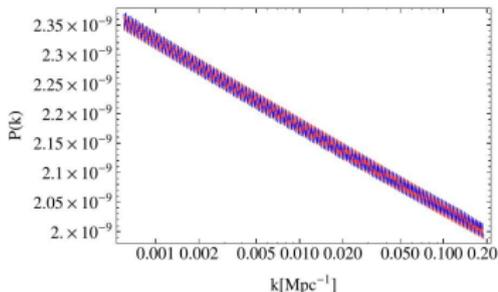
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Local features and large scale suppression

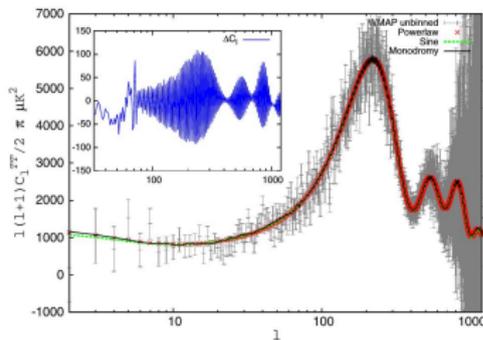
Vilenkin, Ford 1982; Powell, Kinney 2006



Non-local features : axion monodromy model

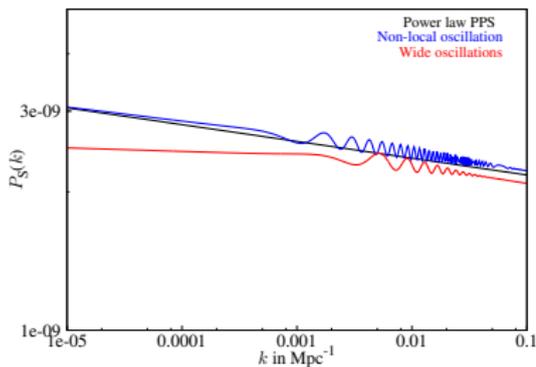
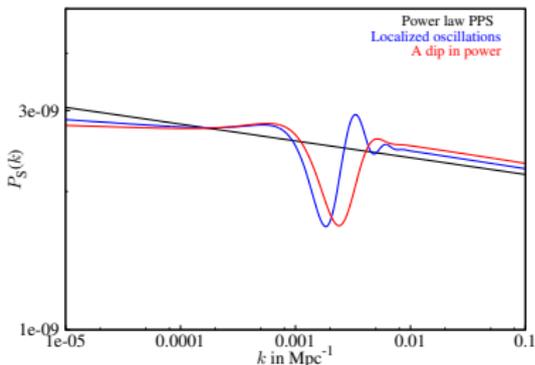
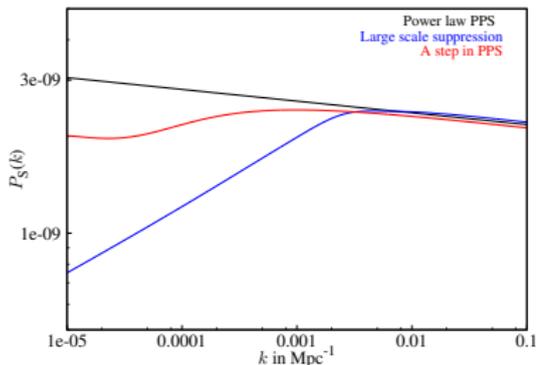


Flauger et al. 2009, 2014

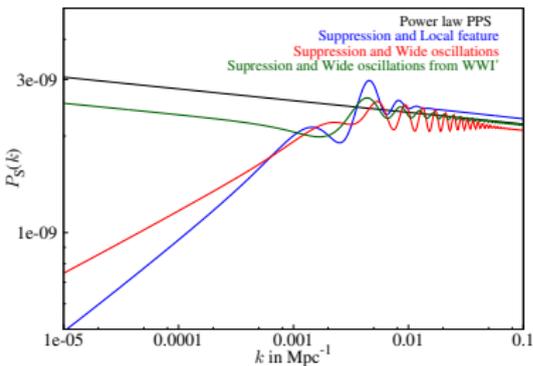
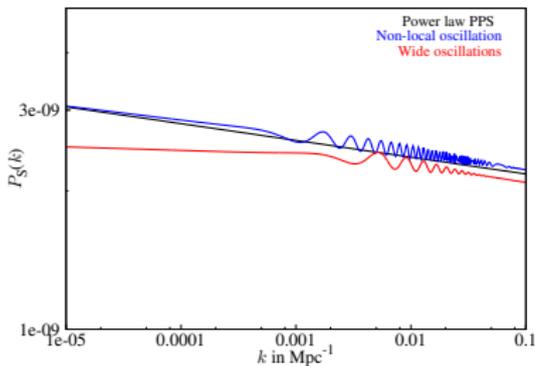
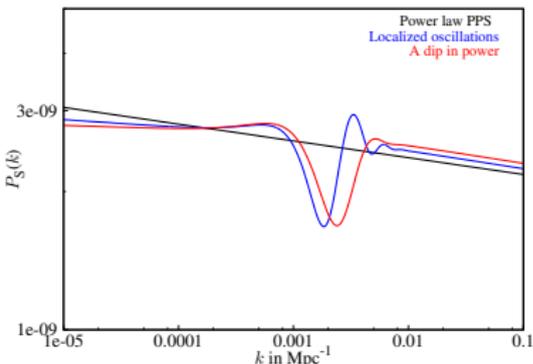
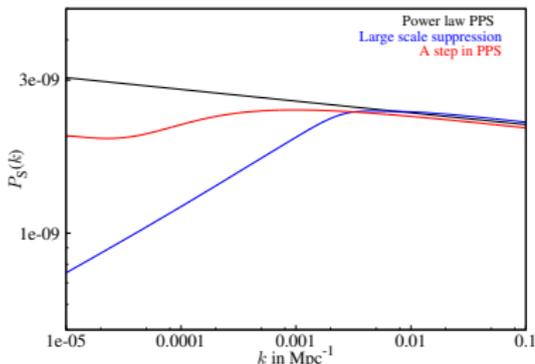


Aich, Hazra, Sriramkumar, Souradeep, [PRD 2013](#),
Meerburg et. al. 2013 - 2014

Classes of features in WWI



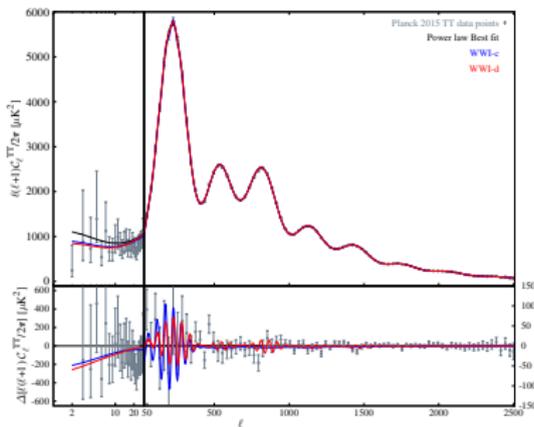
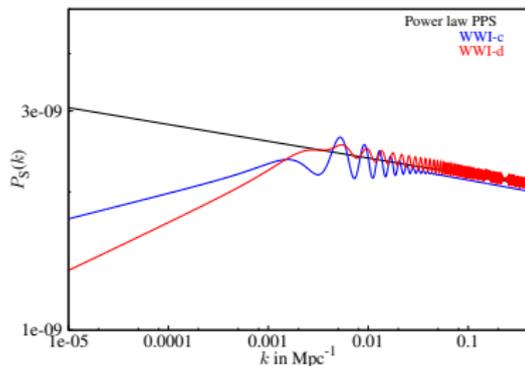
Classes of features in WWI



Wiggly Whipped Inflation : Planck 2015



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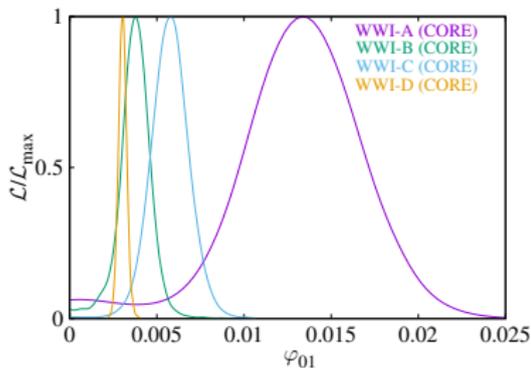
WWI-c provides ~ 10 improvement in χ^2 fit. This improvement comes from low- ℓ TEB and high- ℓ E data.

WWI-d provides $\sim 12 - 14$ improvement while most of the improvement comes owing to the inability of standard model in fitting the temperature and polarization datasets in a combination.

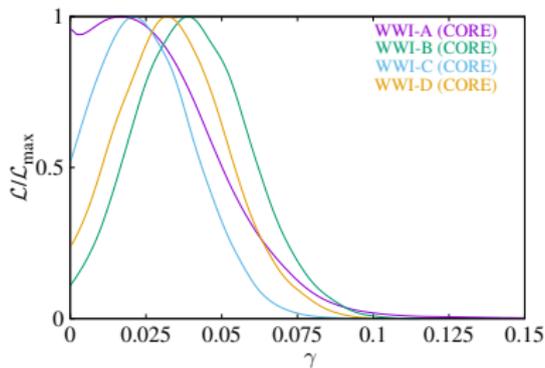
Hazra, Shafieloo, Smoot, Starobinsky, [JCAP 2016](#)

With Cosmic Origins Explorer (CORE)-like survey specification

Wiggles



Suppression



- Large scale suppressions can not be detected with high significance
- Some of the intermediate and small scale oscillations can be detected, if present

Hazra, Paoletti, Ballardini, Finelli, Shafieloo, Smoot, Starobinsky, **JCAP 2018**



Reconstructing reionization history

How the Universe got reionized



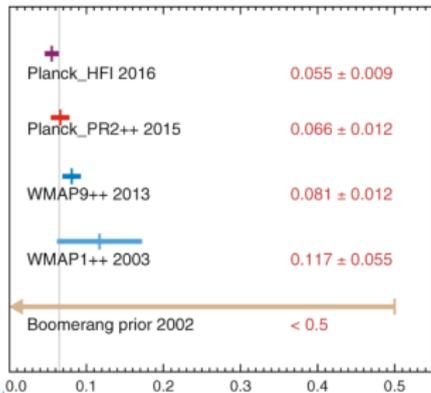
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The common practice is to parametrize the ionized hydrogen fraction/free electron fraction as a function of redshift

In CMB study usually a Tanh model of nearly-instantaneous reionization has been used [Planck has used some asymmetric models as well]

Solving the ionization equation by parametrizing ionizing photon emission and recombination time is another way to model the history

The integrated optical depth constraint that we have till now:



$$\text{Planck 2018} = 0.054 \pm 0.007$$

Parametrizing ionized hydrogen



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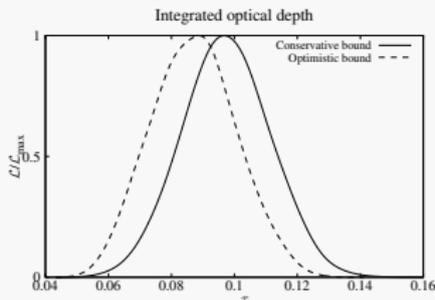
Introduced **Poly-Reion** (a way to construct the history of reionization) and explored possible ways the Universe could have been reionized.

Compared to the conventional nearly instantaneous reionization history we find Planck-2015 prefers an extended history, **higher optical depth**

First free form history of reionization respecting physical bounds on the free electron fraction

Constraints on the optical depth

$$\tau = \int \sigma_T n_e(z) dl$$



Parametrizing ionized hydrogen



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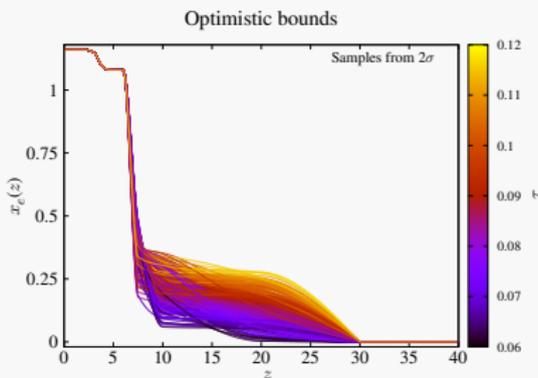
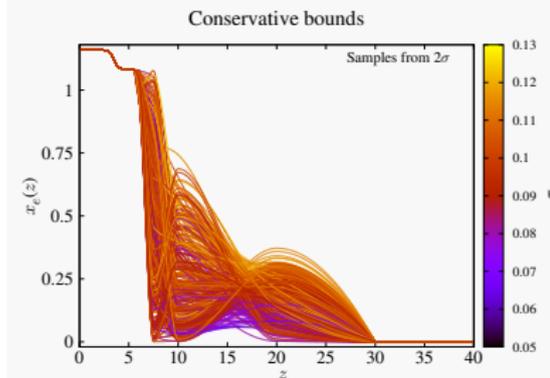
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Only CMB

CMB+low redshift observables



Hazra and Smoot, JCAP 2017

Parametrizing ionized hydrogen



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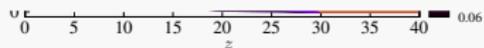
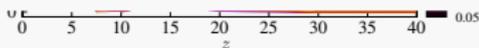
Conservative bounds



Optimistic bounds



Poly-reion formalism has been used by different groups, Villanueva-Domingo *et. al.* 2018; Millea and Bouchet 2018 (in a modified form) and Planck 2018: cosmological parameters



Hazra and Smoot, JCAP 2017



- Why free form solution is important:
 - Not biased towards particular phenomenological construct
 - Solving ionization equation provides realistic ionization fractions
 - Constraints are conservative
 - Physical parameters can be obtained in post-process
- Why complete CMB data is needed:
 - Correlations with other parameters are important
 - Use of τ -only constraints can be biased towards the baseline model
 - Analysis with adding other dataset will be robust

Hazra, Paoletti, Finelli, Smoot, [arXiv 2019](#)

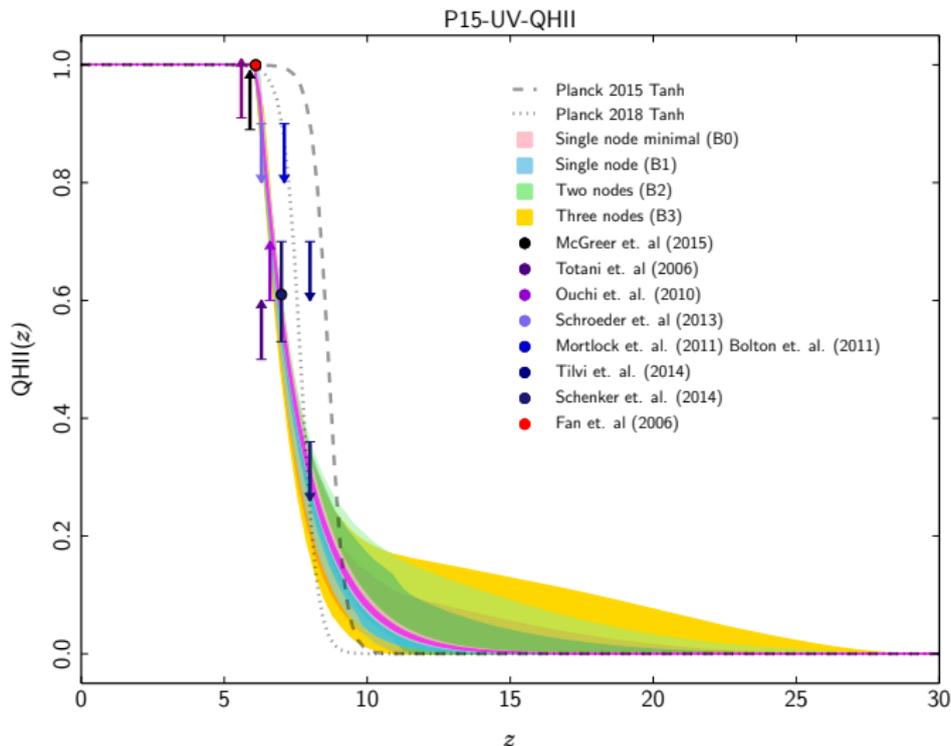
- 1 **Planck 2015 angular power spectrum and lensing (P15)**
- 2 **UV luminosity density from Hubble Frontier Fields (HFF) data (UV)**
Spanning $z \sim 6 - 11$, we use Bouwens et. al. 2015 and Ishigaki et. al. 2018
- 3 **Constraints from Lyman- α (QHII)**
Spanning $z \sim 6 - 8$, we use Fan et. al. 2006, Schroeder et. al. 2013, Schenker et. al. 2014
Using parametric form and Planck-2016 τ and UV luminosity density from HFF, recent works are Gorce et. al. 2018, Ishigaki et. al. 2018; a new analysis with free form source function appeared recently Mason et. al. 2019

Hazra, Paoletti, Finelli, Smoot, [arXiv 2019](#)

Free form solution: constraints



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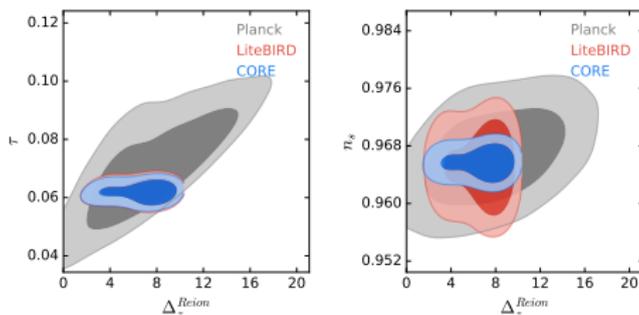


Degeneracies and how to break them



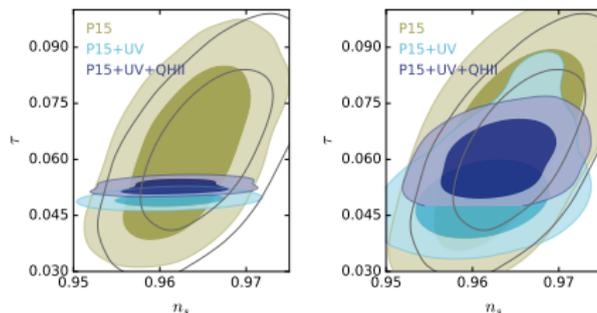
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Wait for better data:



Hazra, Paoletti, Finelli, Smoot, [JCAP 2018](#)

Use astrophysical observations



Hazra, Paoletti, Finelli, Smoot, [arXiv 2019](#)



- Analysis of inflation and reionization with Planck 2018 data (released in July 2019)
- Mock reconstruction with Euclid
- Joint forecast of features with Euclid
- Breaking correlations between inflation and reionization, chances with JWST.

Thank you